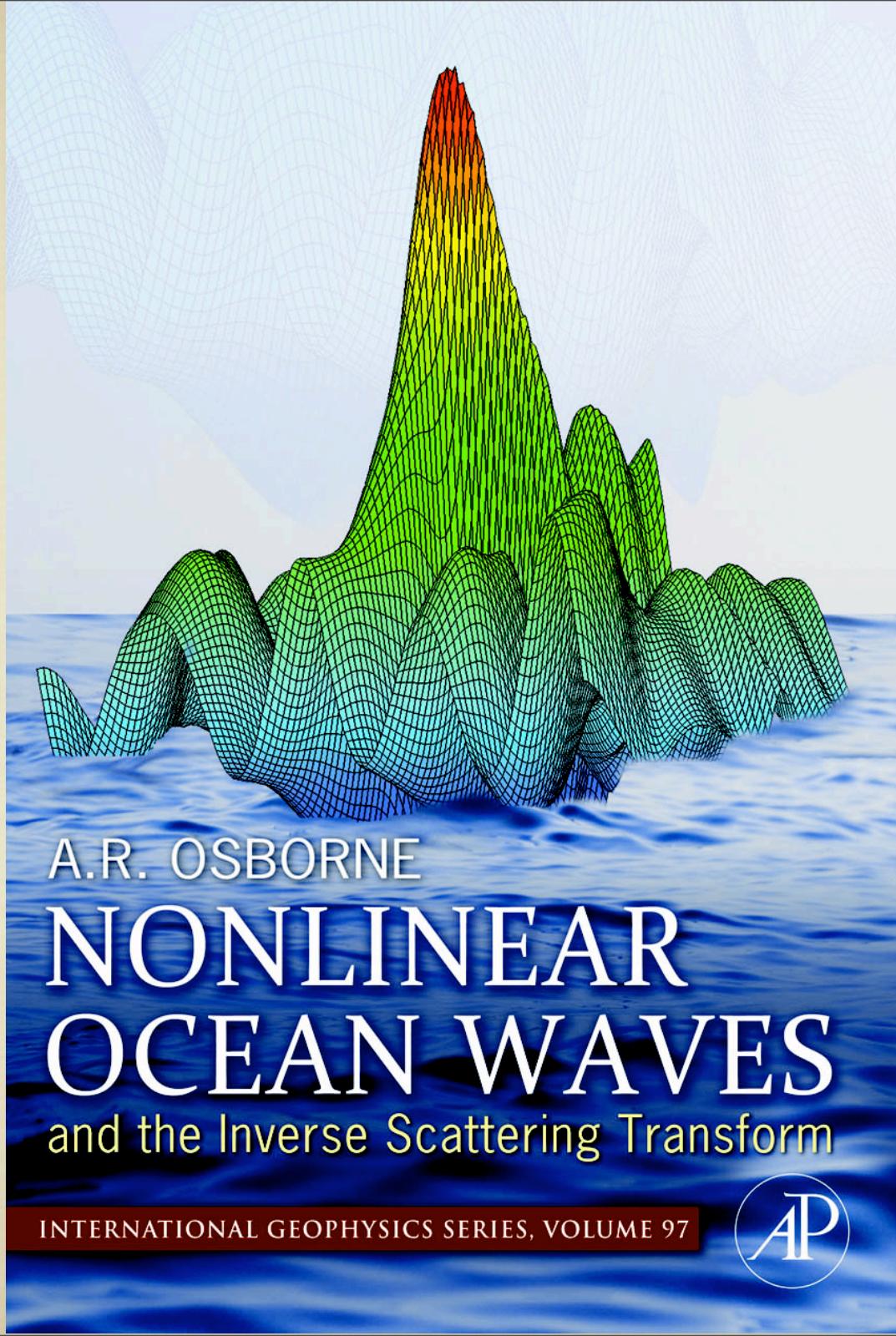


# *Modeling Waves in Shallow (and Deep) Water*

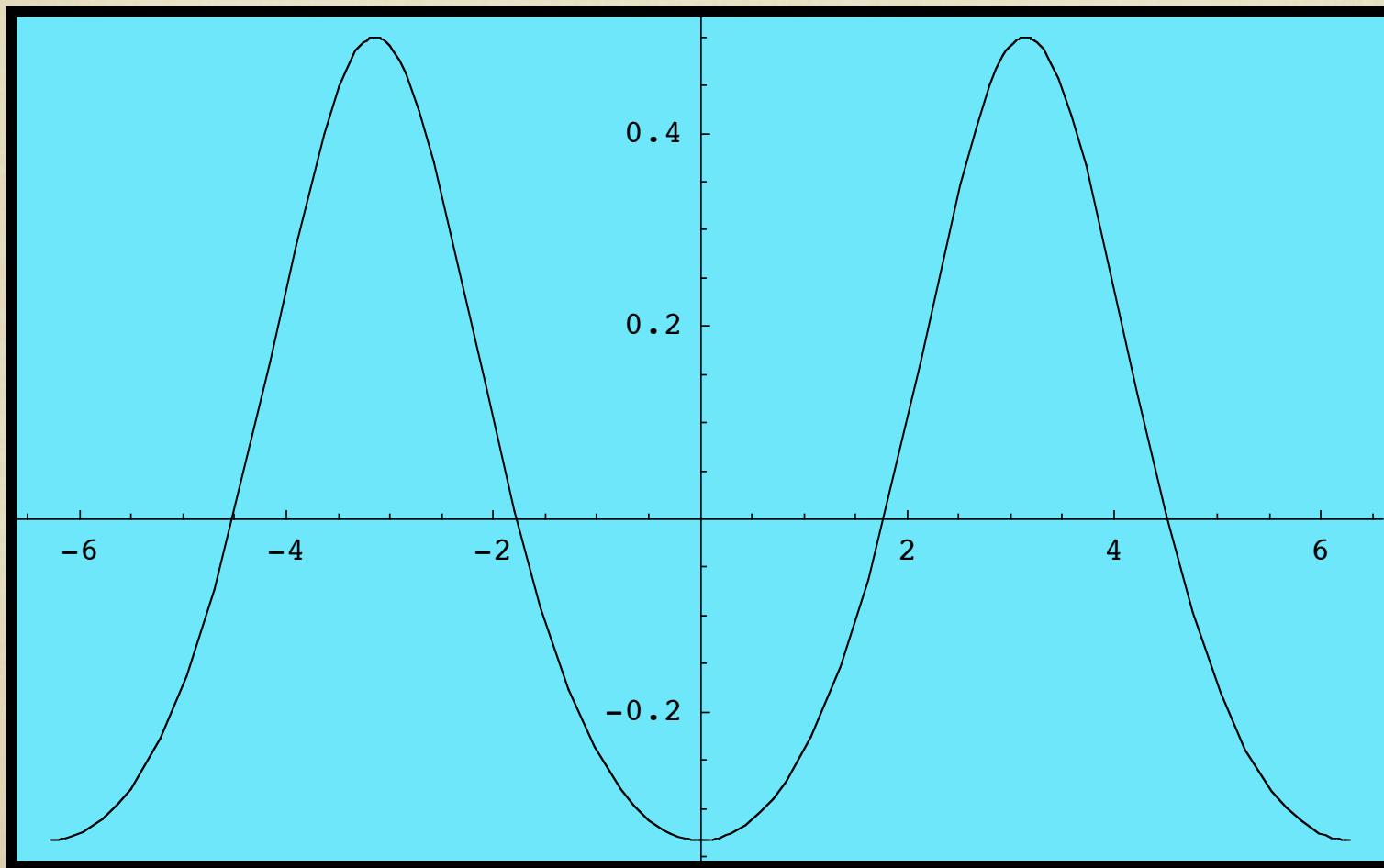
*Al Osborne  
Don Resio*

**FUNDERS: ONR AND ARMY CORP OF ENGINEERS**

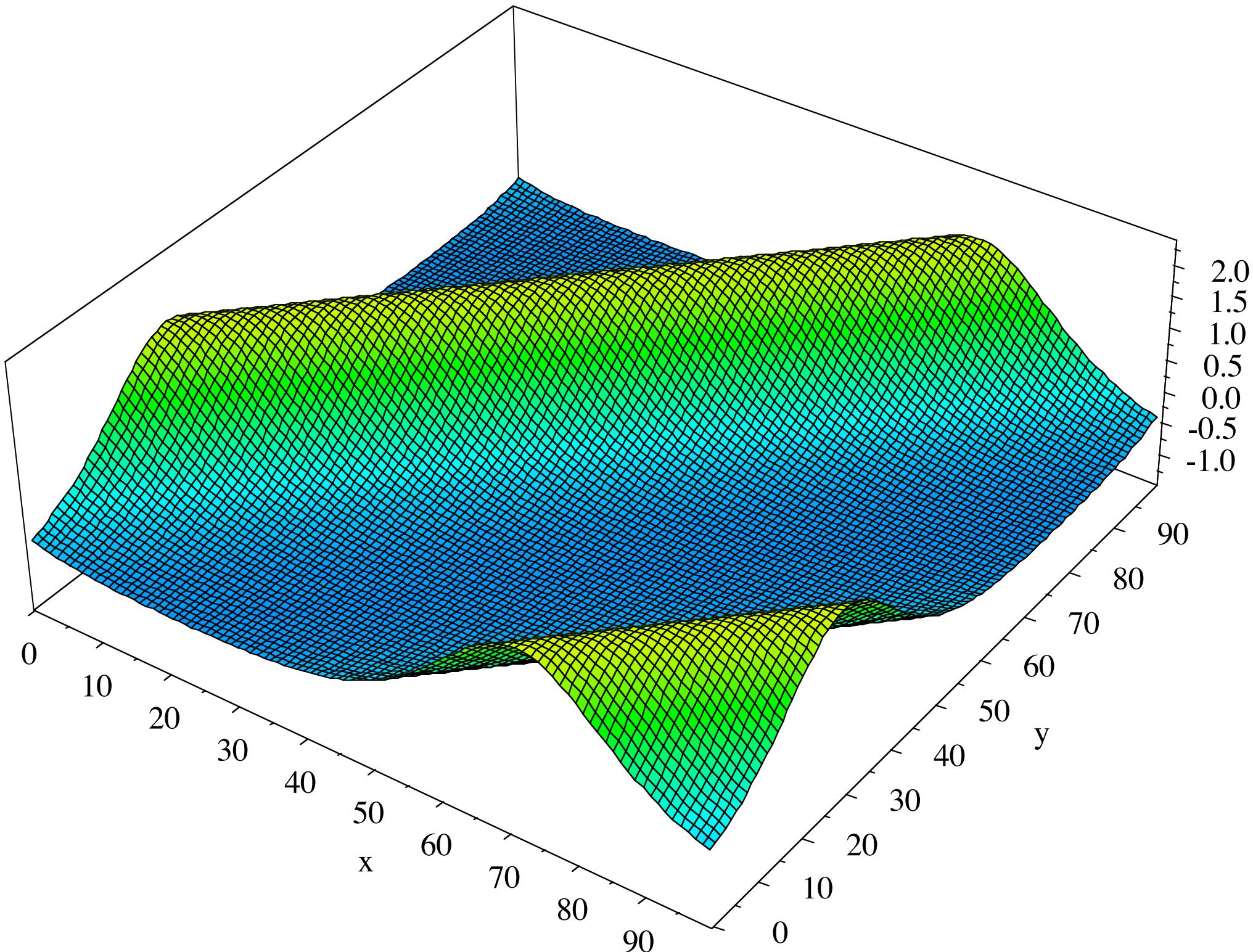
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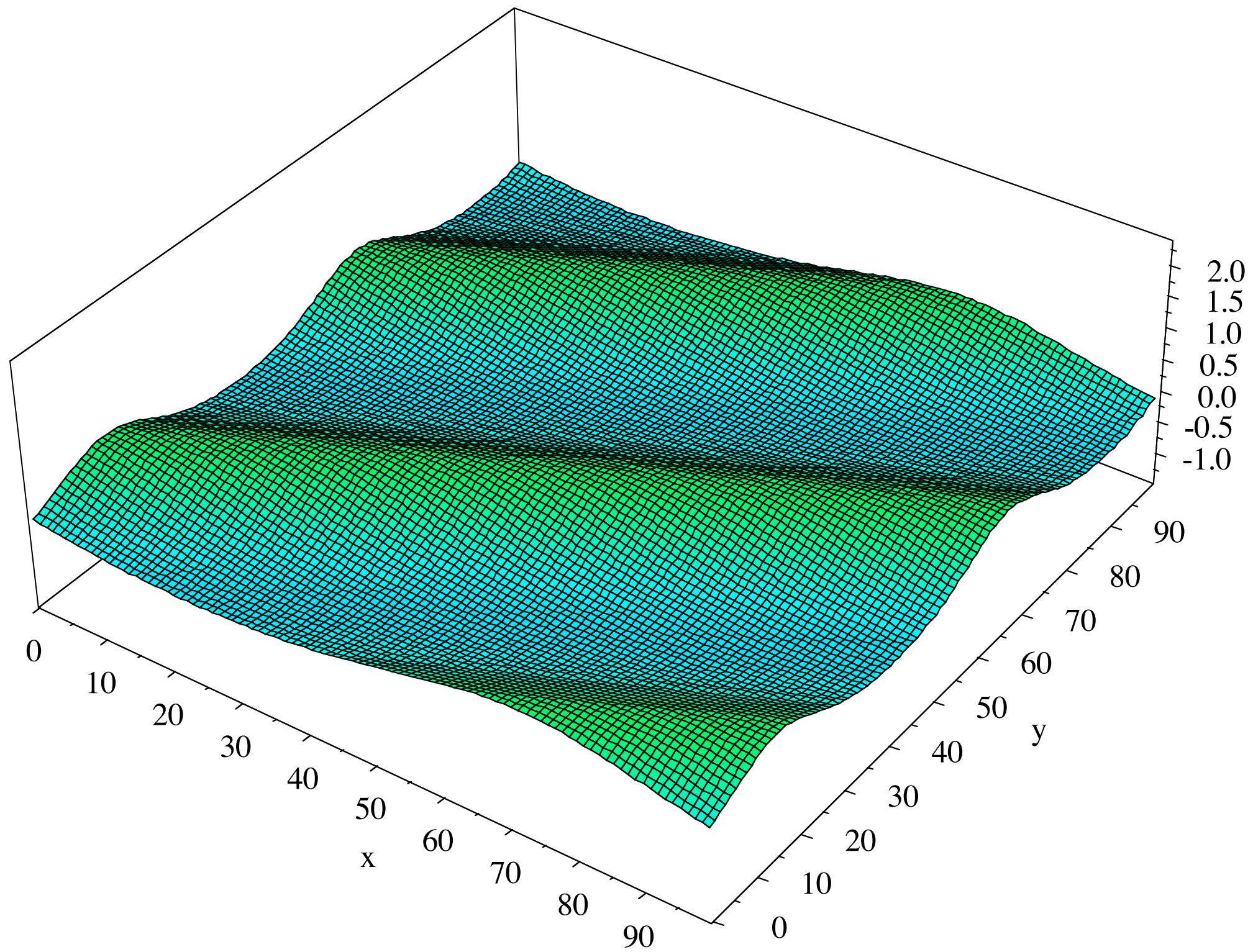


# *Stokes Wave Nonlinearity*

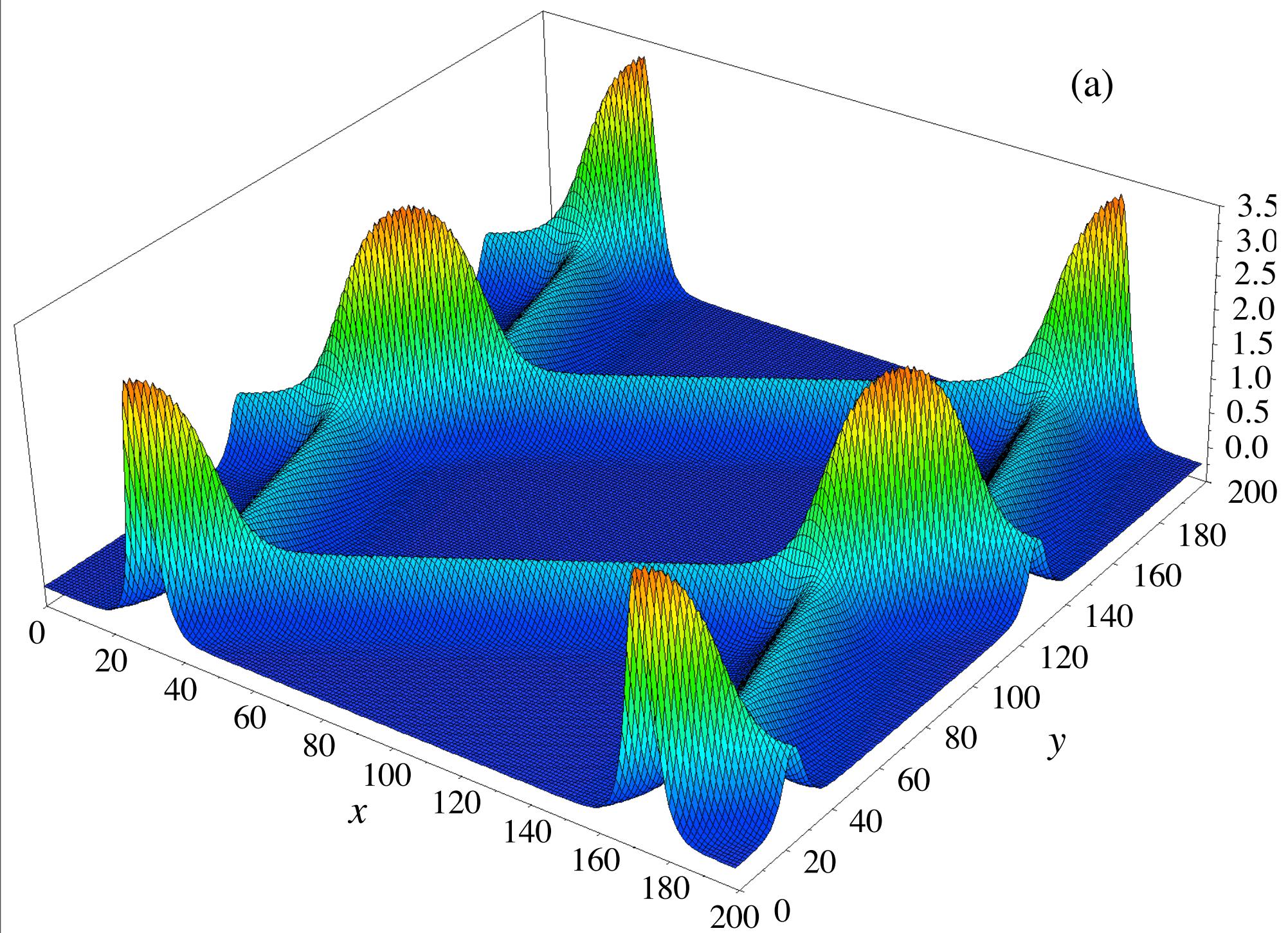


# Nonlinear Interactions for Directional Stokes Waves





(a)





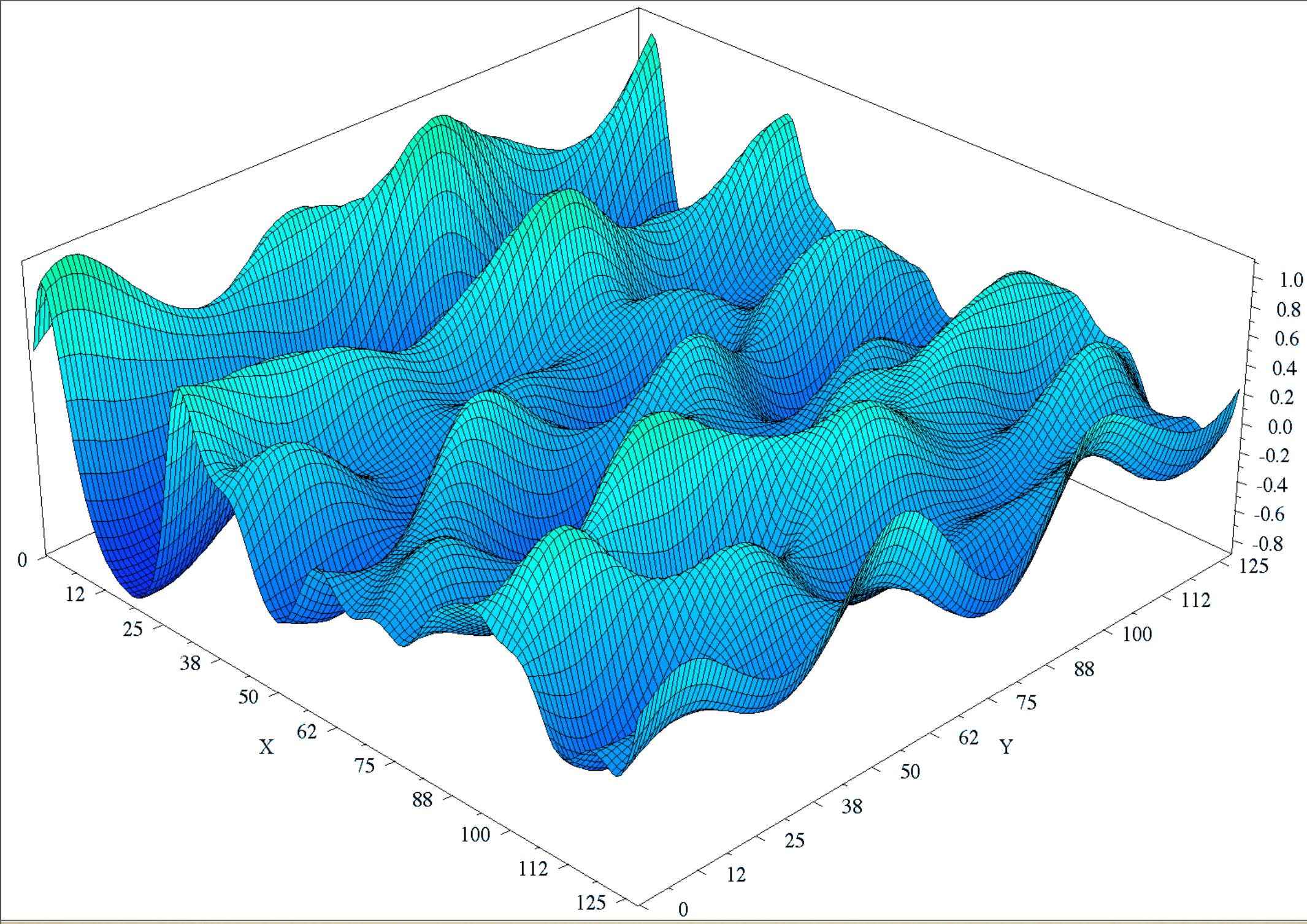
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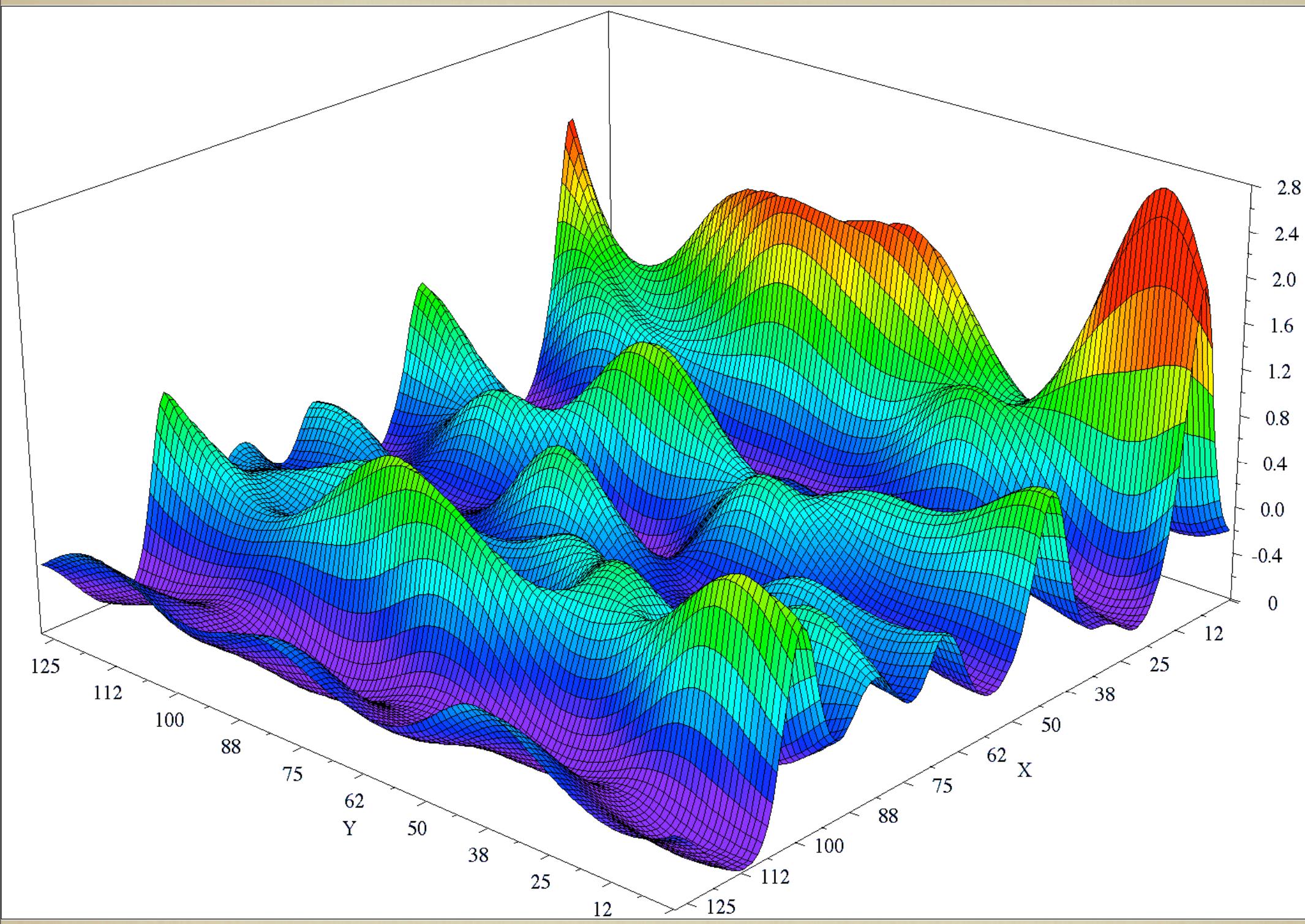
## *Kadomtsev and Petviashvili (KP) equation*

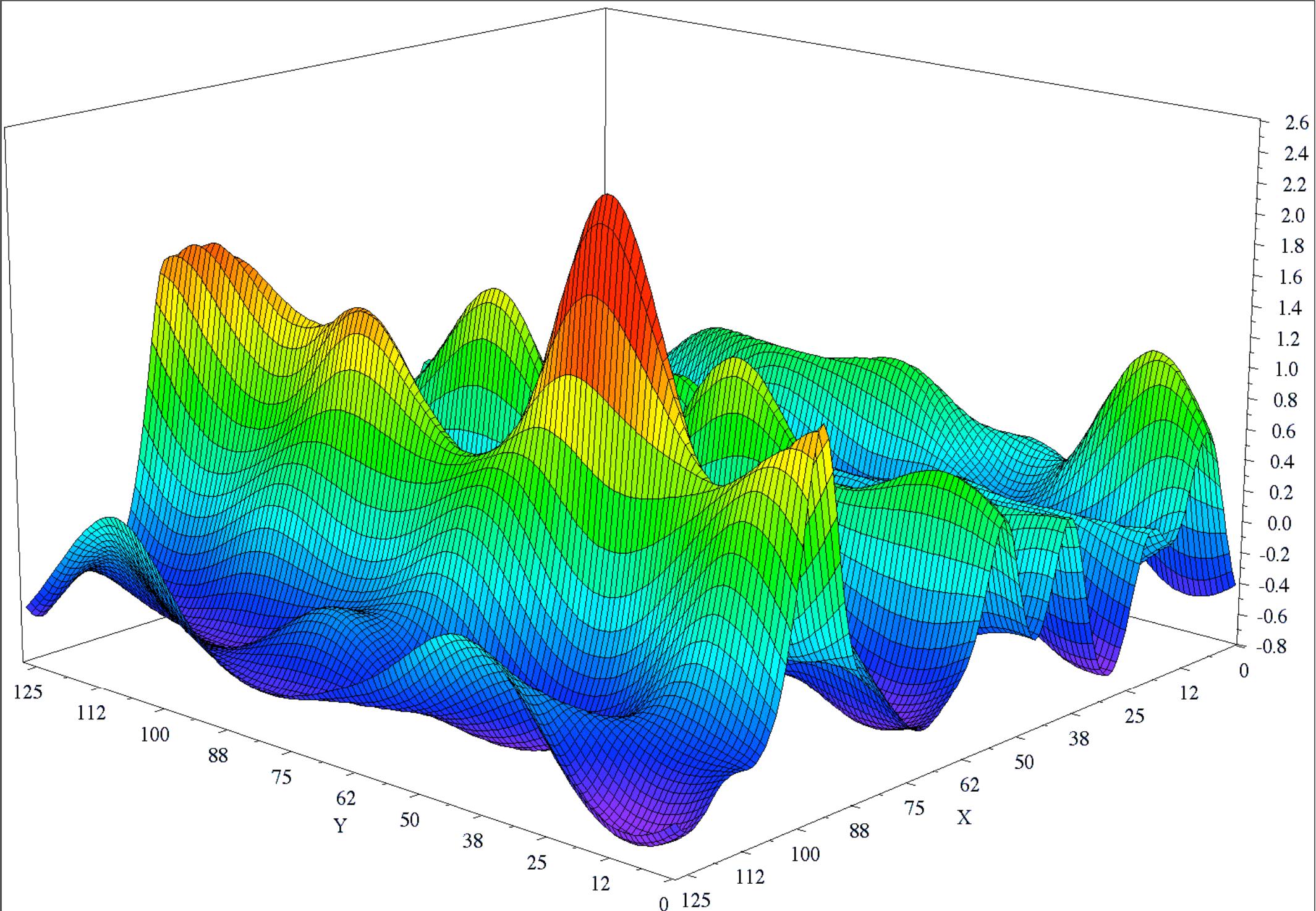
$$\eta_t + c_o \eta_x + \alpha \eta \eta_x + \beta \eta_{xxx} + \gamma \partial_x^{-1} \eta_{yy} = 0$$

$$\eta(x, y, t) = 2 \frac{\partial^2}{\partial x^2} \ln \theta(x, y, t)$$

$$\theta(x, y, t) = \sum_{\mathbf{m} \in \mathbf{Z}} e^{-\frac{1}{2} \mathbf{m} \cdot \mathbf{B} \mathbf{m} + i \mathbf{m} \cdot \mathbf{k} x + i \mathbf{m} \cdot \mathbf{l} y - i \mathbf{m} \cdot \boldsymbol{\omega} t + \mathbf{m} \cdot \boldsymbol{\phi}}$$







# Modulational Instability

# DAVEY-STEWARTSON EQUATIONS

$$i\Psi_\tau + \lambda\Psi_{XX} + \mu\Psi_{YY} + \chi|\Psi|^2\Psi = \chi_o\Psi\Phi_X$$

$$\Phi_{XX} + \Phi_{YY} = -\beta\left(|\Psi|^2\right)_X$$

$$\eta(x,y,t) \approx \frac{i\omega\sqrt{gk_o}}{gk_o^2}\Psi(x,y,t)e^{ik_ox-i\omega_ot} + c.c. + ...$$

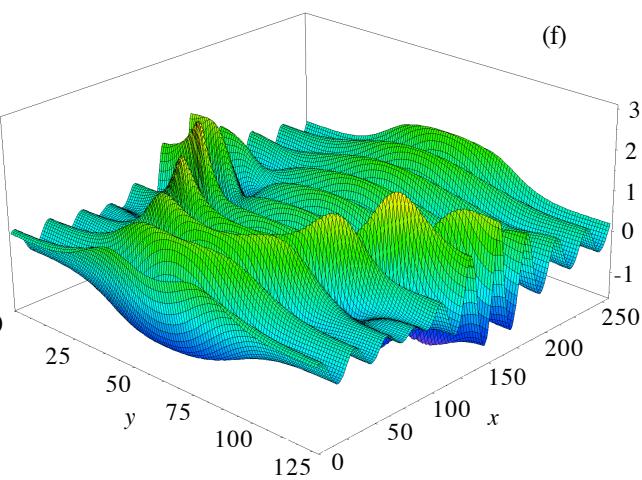
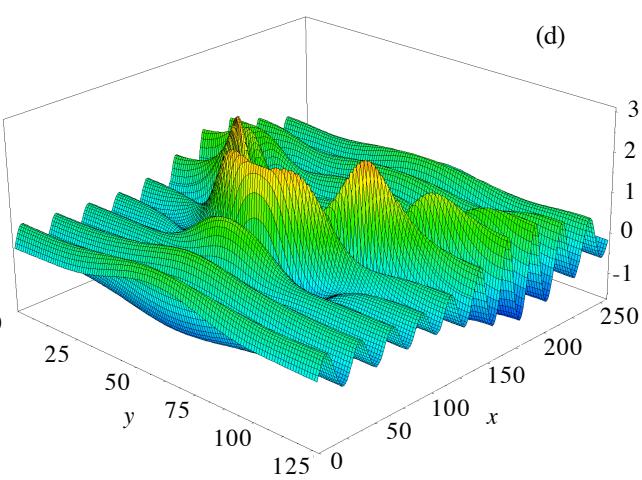
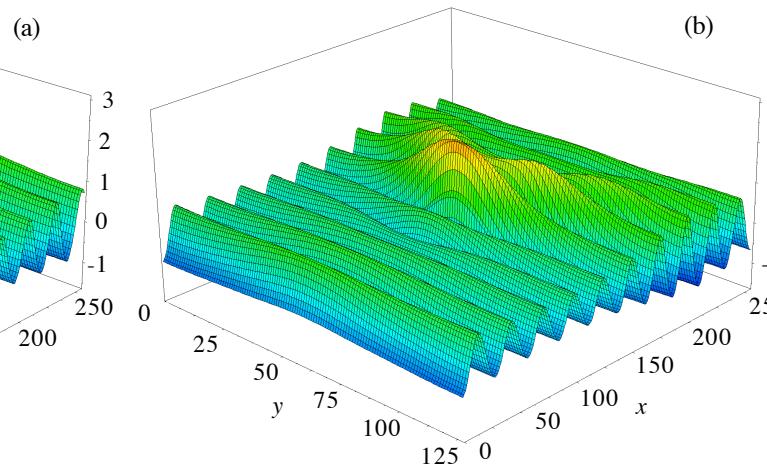
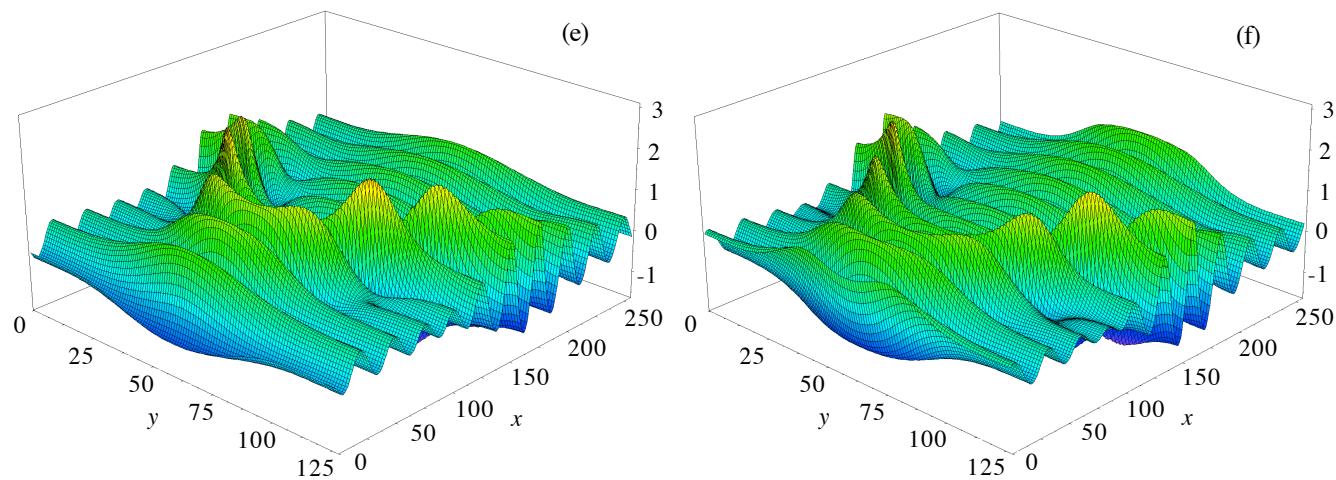
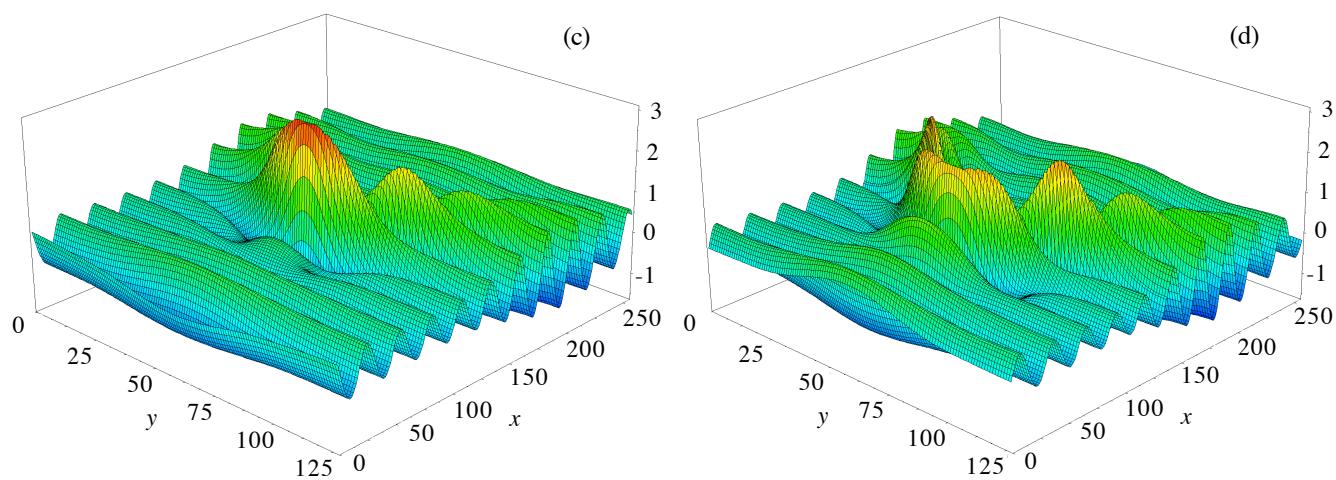
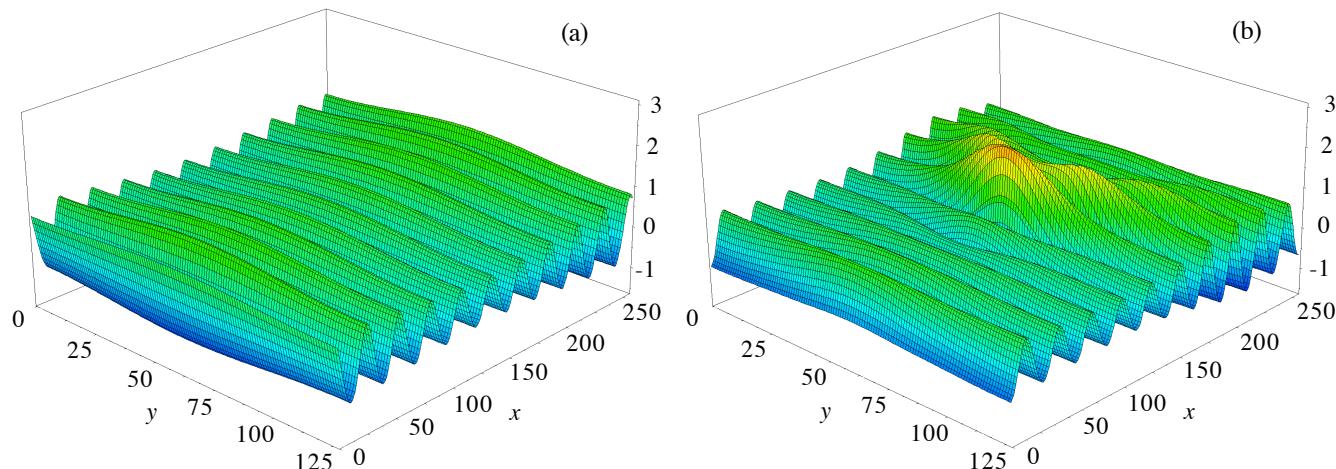
$$\phi(x,y,t) \approx \sqrt{\frac{g}{k^3}}\left(\Phi(x,y,t) + \frac{\cosh k(z+h)}{\cosh kh}\Psi(x,y,t)e^{ikx-i\omega t} + c.c.\right)$$

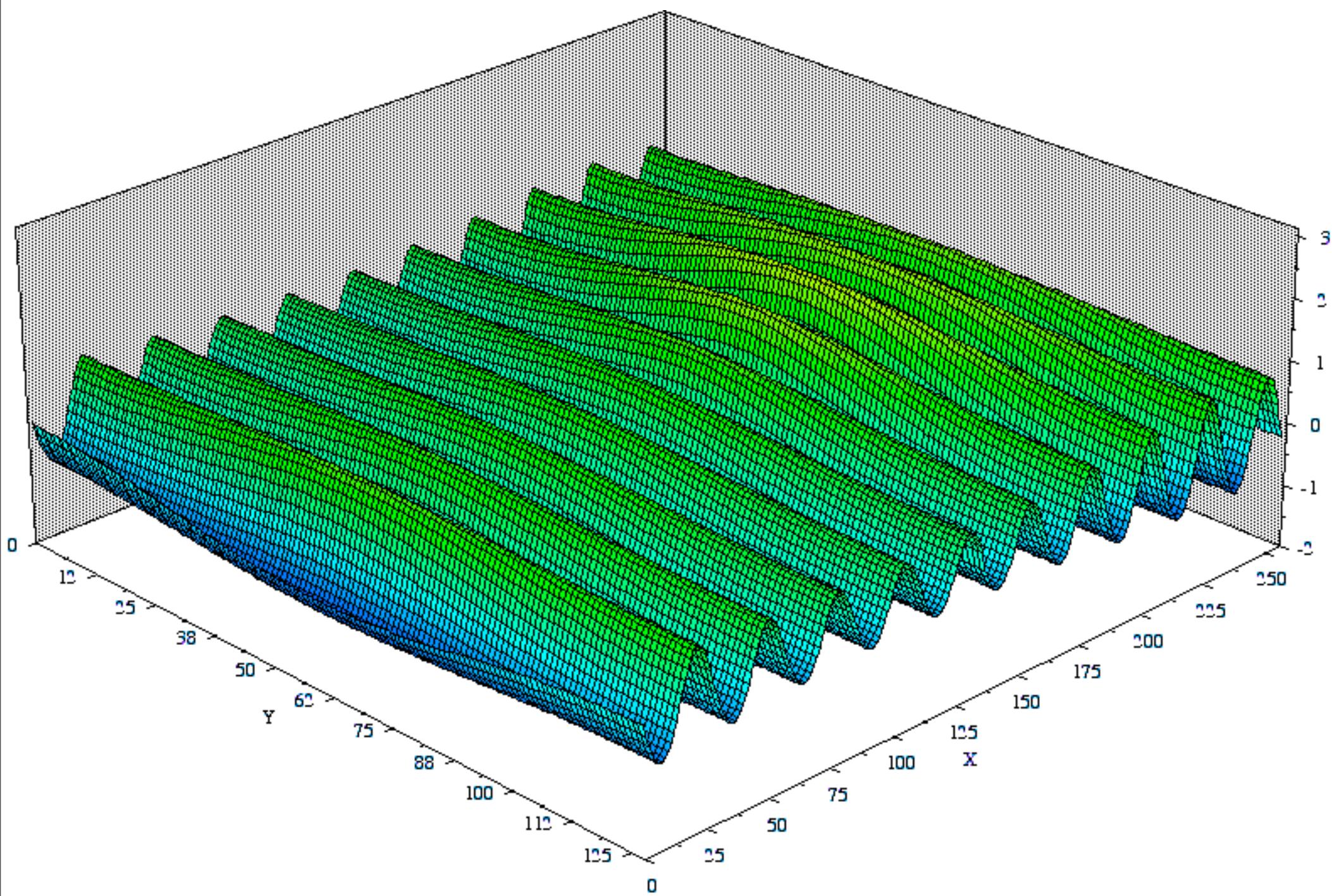
# Instabilities in Deep Water Wave Trains

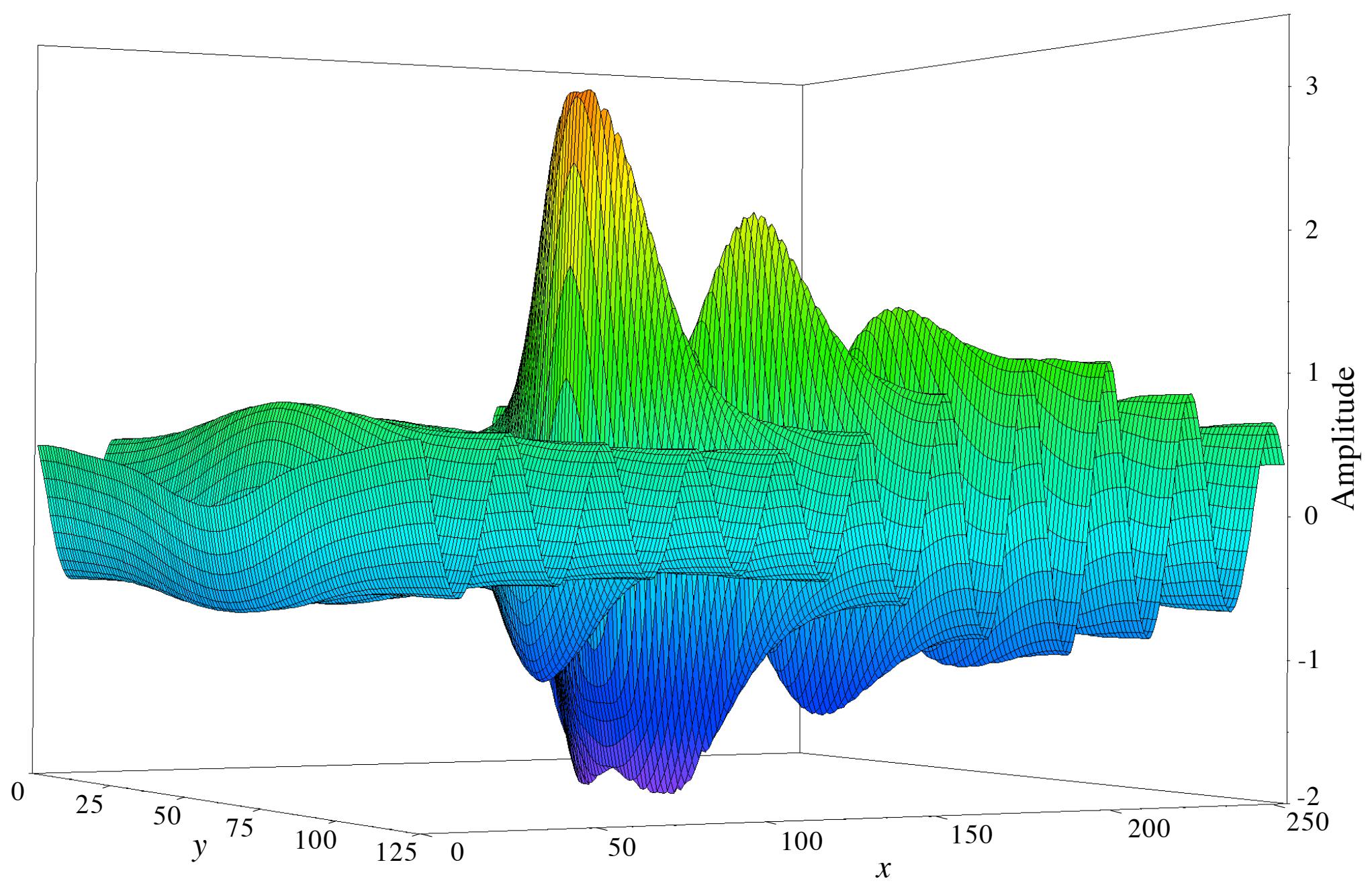
Deep water wave trains have *two kinds of waves* (at leading order):

- (1) *Stokes waves* - quasi linear wave trains described by a spectrum.
- (2) *Unstable wave packets* governed by Benjamin-Fier instability.

**Threshold** - Only *Stokes waves* occur below the BF instability threshold. When the waves get big enough, rapidly enough *unstable packets* begin to form. The extreme case is called a *rogue sea*.







## Conclusions:

Nonlinear Fourier Analysis:

The Inverse Scattering Transform for Periodic Boundary Conditions

- (1) Physics of nonlinear waves as a nonlinear Fourier decomposition of nonlinear wave trains in terms of nonlinear Fourier components, Including **sine waves, Stokes waves, solitons, holes, unstable (rogue) Modes, shock waves, tabletop solitons, vortices, etc.**
- (2) IST gives us approaches for **analyzing directional data** to obtain the nonlinear Fourier spectrum
- (3) Hyperfast numerical methods based upon IST are **three orders of magnitude faster** than FFT approaches.
- (4) Perfectly parallel algorithms give us many more orders of magnitude speed up un computers with thousands of processors.

# *Deterministic Waves: How Linear Fourier Analysis Works*

Given a linear wave equation (shallow water):

$$u_t + c_o u_x + \beta u_{xxx} = 0$$

$$c_o = \sqrt{gh}$$

$$\beta = c_o h^2 / 6$$

Fourier Transform:

$$u_n(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} u(x, t) e^{-ikx} dx$$

Time dependence of Fourier components

$$u_n(t) = u_n(0) e^{-i\omega t}$$

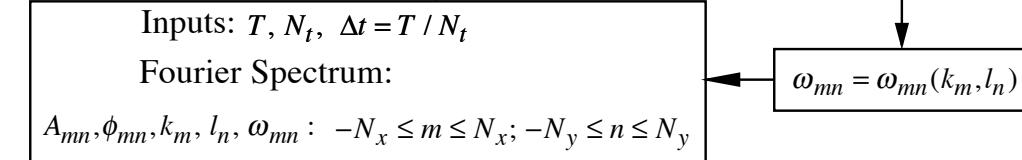
Dispersion Relation

$$\omega = \omega(k) = c_o k - \beta k^3$$

Cauchy Problem:

$$u(x, t) = \sum_{n=-\infty}^{\infty} u_n(0) e^{i(k_n x - \omega_n t)}$$

# Solving Linear PDEs with the Fourier Transform



Dispersion Relation

$$\omega_{mn} = \omega_{mn}(k_m, l_n)$$

Preprocessor

$$t = 0$$

yes

$$t > T ?$$

Time varying  
Fourier Coefficients

$$\eta_{mn}(t) = \eta(k_m, l_n, t) = A_{mn} e^{-i\omega_{mn}t + i\phi_{mn}}; -N_x \leq m \leq N_x; -N_y \leq n \leq N_y$$

$$t = t + \Delta t$$

Solution of Linear PDE

yes

$$t > T ?$$

FFT Solution  
of Linear PDE

$$\eta(x, y, t) = \sum_{m=-N_x/2}^{N_x/2} \sum_{n=-N_y/2}^{N_y/2} \eta_{mn}(t) e^{ik_m x + il_n y}$$

$$t = t + \Delta t$$

End

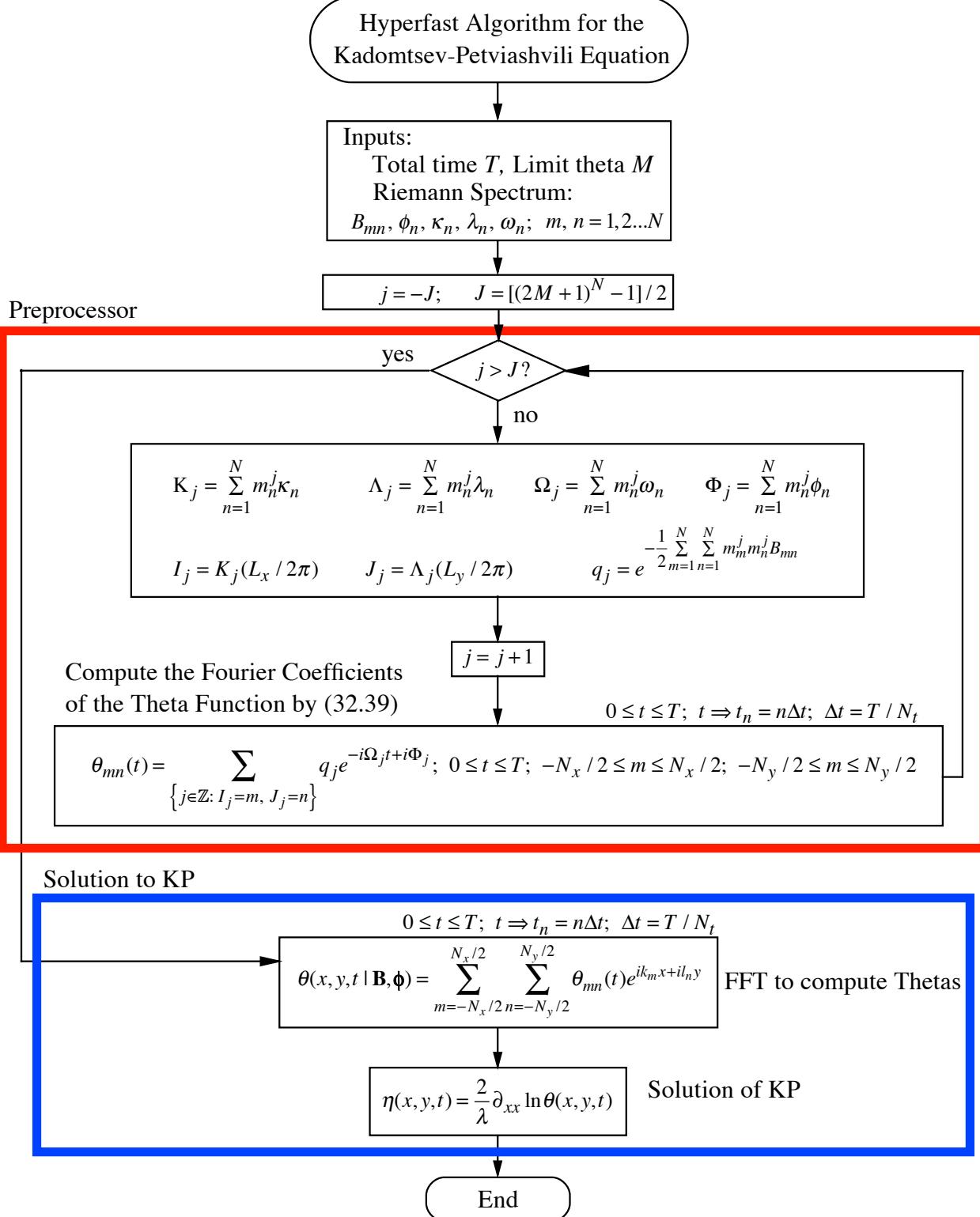
*Kadomtsev and Petviashvili (KP) equation*

## *Kadomtsev and Petviashvili (KP) equation*

$$\eta_t + c_o \eta_x + \alpha \eta \eta_x + \beta \eta_{xxx} + \gamma \partial_x^{-1} \eta_{yy} = 0$$

$$\eta(x, y, t) = 2 \frac{\partial^2}{\partial x^2} \ln \theta(x, y, t)$$

$$\theta(x, y, t) = \sum_{\mathbf{m} \in \mathbf{Z}} e^{-\frac{1}{2} \mathbf{m} \cdot \mathbf{B} \mathbf{m} + i \mathbf{m} \cdot \mathbf{k} x + i \mathbf{m} \cdot \mathbf{l} y - i \mathbf{m} \cdot \boldsymbol{\omega} t + \mathbf{m} \cdot \boldsymbol{\phi}}$$



## *Partial Theta Summations*

$$\theta(x, y, t) = \sum_{m_1=-M}^M \sum_{m_2=-M}^M \dots \sum_{m_N=-M}^M e^{-\frac{1}{2}\mathbf{m} \cdot \mathbf{B}\mathbf{m} + i\mathbf{m} \cdot \mathbf{k}x + i\mathbf{m} \cdot \mathbf{l}y - i\mathbf{m} \cdot \boldsymbol{\omega}t + \mathbf{m} \cdot \boldsymbol{\phi}}$$

$$\mathbf{m} = [m_1, m_2 \dots m_N]$$

*Exponential*

$$(2M+1)^N \sim 10^N$$

**Theorem:** The Riemann theta function can be reduced to the linear Fourier transform with time dependent coefficients.

$$\theta(x, y, t) = \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} \theta_{mn}(t) e^{ik_m x + il_n y}$$

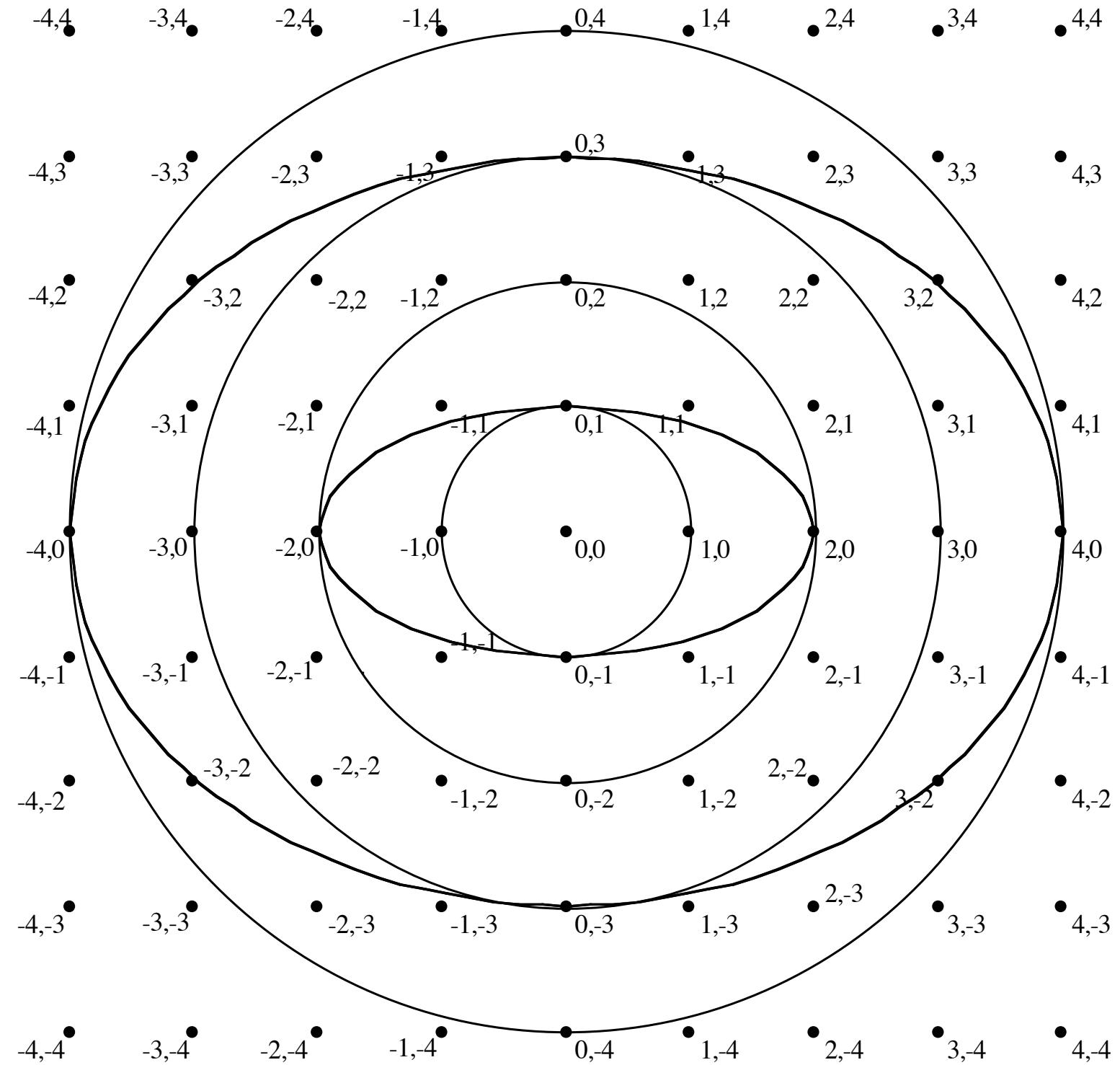
$$\theta_{mn}(t) = \sum_{\substack{\{\mathbf{m} \in \mathbf{Z}: L_x(\mathbf{m} \cdot \mathbf{k})/2\pi = m\} \\ \{\mathbf{m} \in \mathbf{Z}: L_y(\mathbf{m} \cdot \mathbf{l})/2\pi = n\}}} Q_{\mathbf{m}} e^{-i\mathbf{m} \cdot \boldsymbol{\omega} t}; \quad Q_{\mathbf{m}} = e^{-\frac{1}{2}\mathbf{m} \cdot \mathbf{B} \mathbf{m} + i\mathbf{m} \cdot \boldsymbol{\phi}}$$

$$\frac{L}{2\pi} \mathbf{m} \cdot \mathbf{k} = \mathbf{m} \cdot \mathbf{n}; \quad \mathbf{n} = [1, 2, 3 \dots N]; \quad \mathbf{k} = [k_1, k_2 \dots k_N]; \quad \mathbf{k} = \frac{2\pi}{L} \mathbf{n}$$

The *linear Fourier Coefficients are Fourier series in time!*

*Linear Fourier instead*

$$\theta_{mn}(t) = \theta_{mn}(0) e^{-i\omega_{mn} t}$$



# Summing Over the $n$ -Sphere

## ***N-Cube***

$$\theta(x, t \mid \mathbf{B}, \phi) = \sum_{m_1=-M}^M \sum_{m_2=-M}^M \dots \sum_{m_N=-M}^M e^{-\frac{1}{2} \sum_{m=1}^N \sum_{n=1}^N m_m m_n B_{mn}} e^{i \sum_{n=1}^N m_n \kappa_n x - i \sum_{n=1}^N m_n \omega_n t + i \sum_{n=1}^N m_n \phi_n}$$

$$m_1^2 = R^2 \quad \text{line (1D)} \quad \Rightarrow m_1 = R = 2M$$

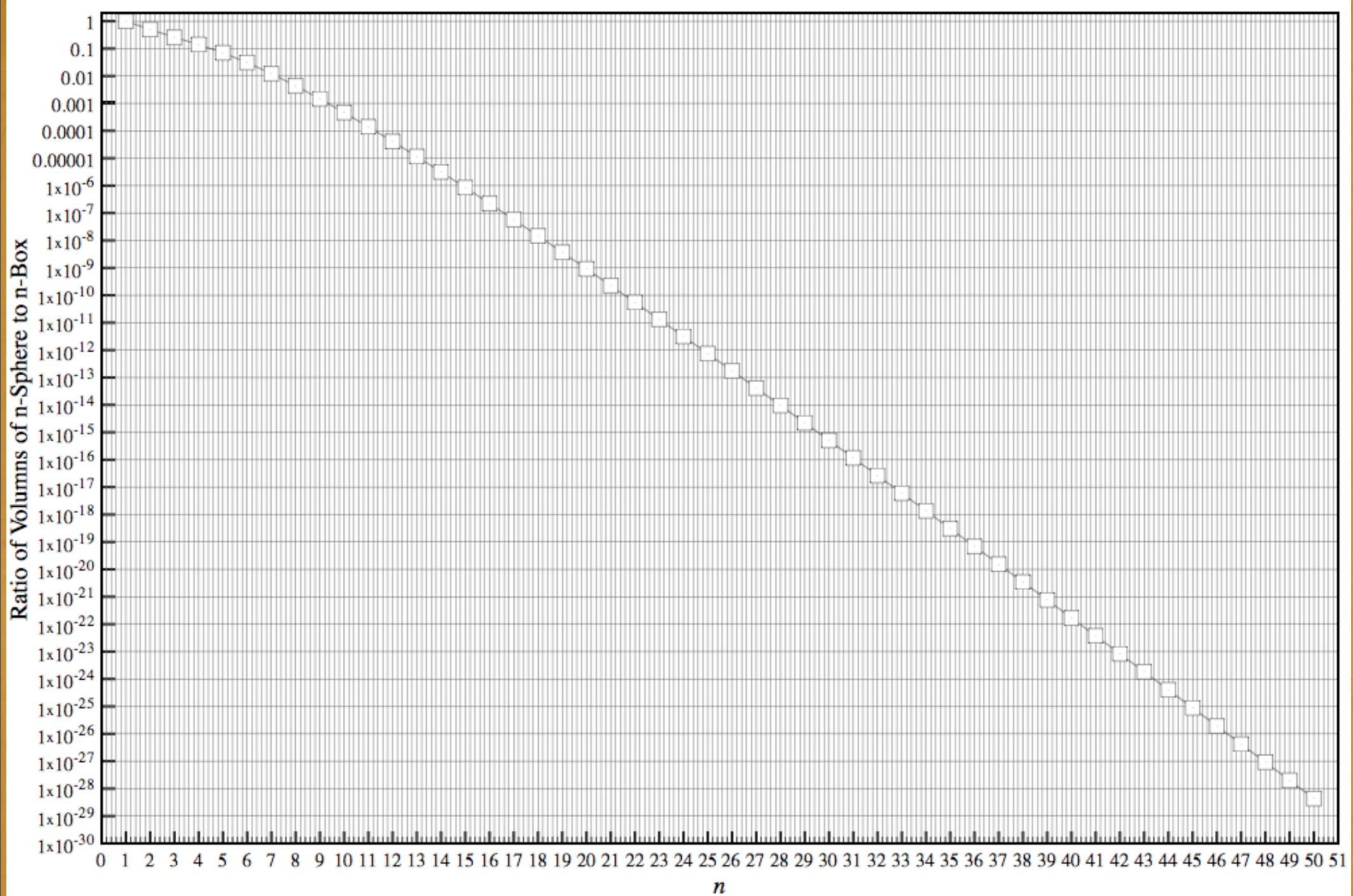
$$m_1^2 + m_2^2 = R^2 \quad \text{circle (2D)} \quad \Rightarrow m_2 = \sqrt{R^2 - m_1^2}$$

$$m_1^2 + m_2^2 + m_3^2 = R^2 \quad \text{sphere (3D)} \quad \Rightarrow m_3 = \sqrt{R^2 - m_1^2 - m_2^2}$$

$$m_1^2 + m_2^2 + m_3^2 + \dots + m_N^2 = R^2 \quad n\text{-sphere (nD)} \quad \Rightarrow m_N = \sqrt{R^2 - m_1^2 - m_2^2 - m_3^2 - \dots}$$

## ***N-Sphere***

$$\theta(x, t) = \sum_{m_1=-M}^M \sum_{m_2=-I\sqrt{M^2-m_1^2}}^{I\sqrt{M^2-m_1^2}} \sum_{m_3=-I\sqrt{M^2-m_1^2-m_2^2}}^{I\sqrt{M^2-m_1^2-m_2^2}} \dots \sum_{m_N=-I\sqrt{M^2-m_1^2-m_2^2-\dots-m_{N-1}^2}}^{I\sqrt{M^2-m_1^2-m_2^2-\dots-m_{N-1}^2}} \times \\ e^{-\frac{1}{2} \sum_{m=1}^N \sum_{n=1}^N m_m m_n B_{mn}} e^{i \sum_{n=1}^N m_n \kappa_n x - i \sum_{n=1}^N m_n \omega_n t + i \sum_{n=1}^N m_n \phi_n}$$



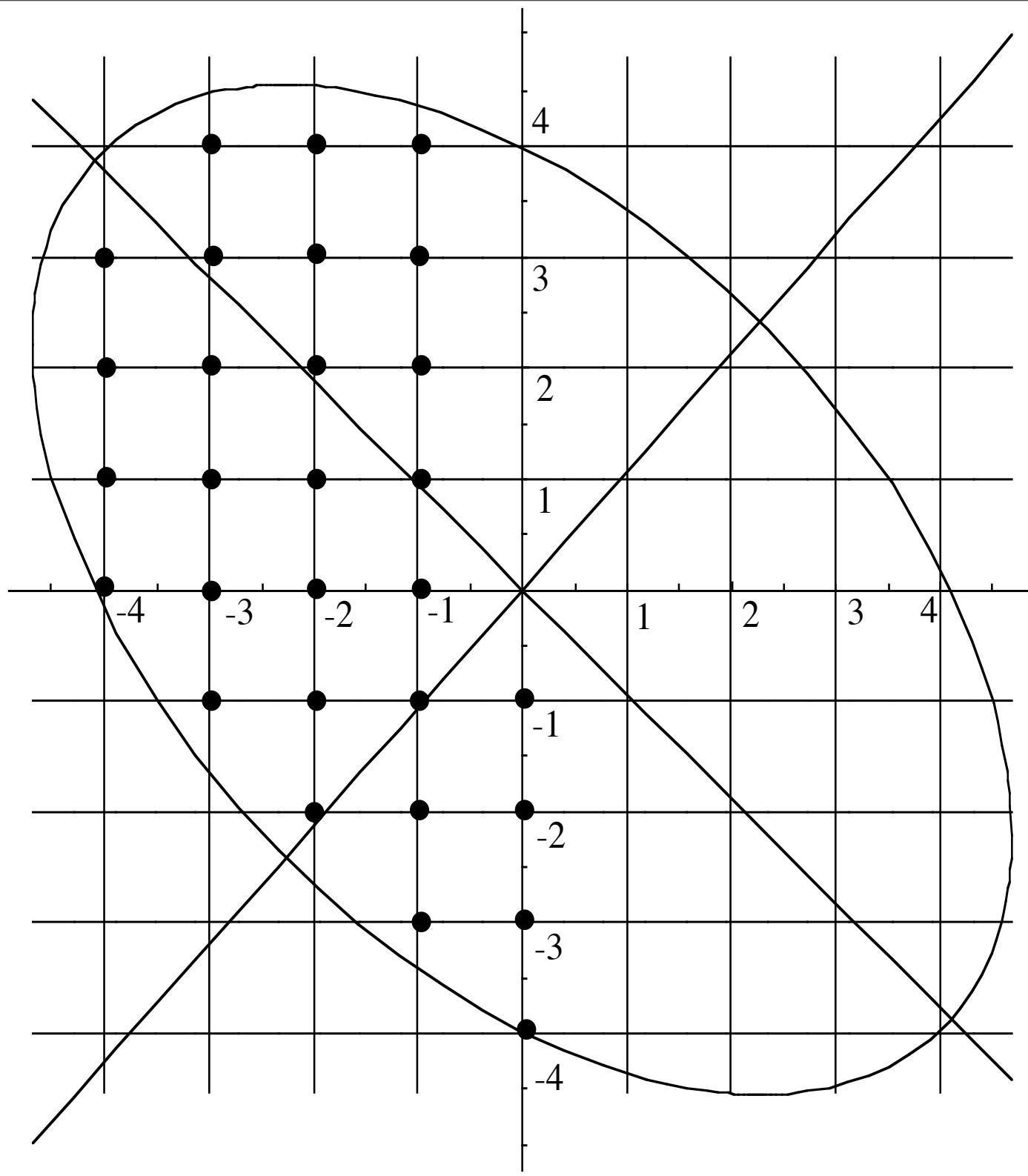
# *Summing Thetas Over the N-Ellipsoid*

$$B_{nn}M_n^2 + 2\left(\sum_{j=1}^{n-1} m_j B_{jn}\right) M_n + \sum_{j=1}^{n-1} \sum_{k=1}^{n-1} m_j m_k B_{jk} + 2 \ln q_o = 0$$

$$M_n^\pm = -\frac{\sum\limits_{j=1}^{n-1} m_j B_{jn}}{B_{nn}} \pm \sqrt{R_{nn}^2 - \frac{\sum\limits_{j=1}^{n-1} \sum\limits_{k=1}^{n-1} m_j m_k B_{jk}}{B_{nn}} + \left(\frac{\sum\limits_{j=1}^{n-1} m_j B_{jn}}{B_{nn}}\right)^2}; \quad R_{nn}^2 = \frac{2|\ln q_o|}{B_{nn}}$$

$$\theta(x,t)=\sum_{m_1=-M_1^-}^{M_1^+}\sum_{m_2=-M_2^-}^{M_2^+}...\sum_{m_N=-M_N^-}^{M_N^+}q_{m_1m_2...m_N}e^{i\sum\limits_{j=1}^Nm_jk_jx-i\sum\limits_{j=1}^Nm_j\omega_jt+i\sum\limits_{j=1}^Nm_j\phi_j}$$

10-16



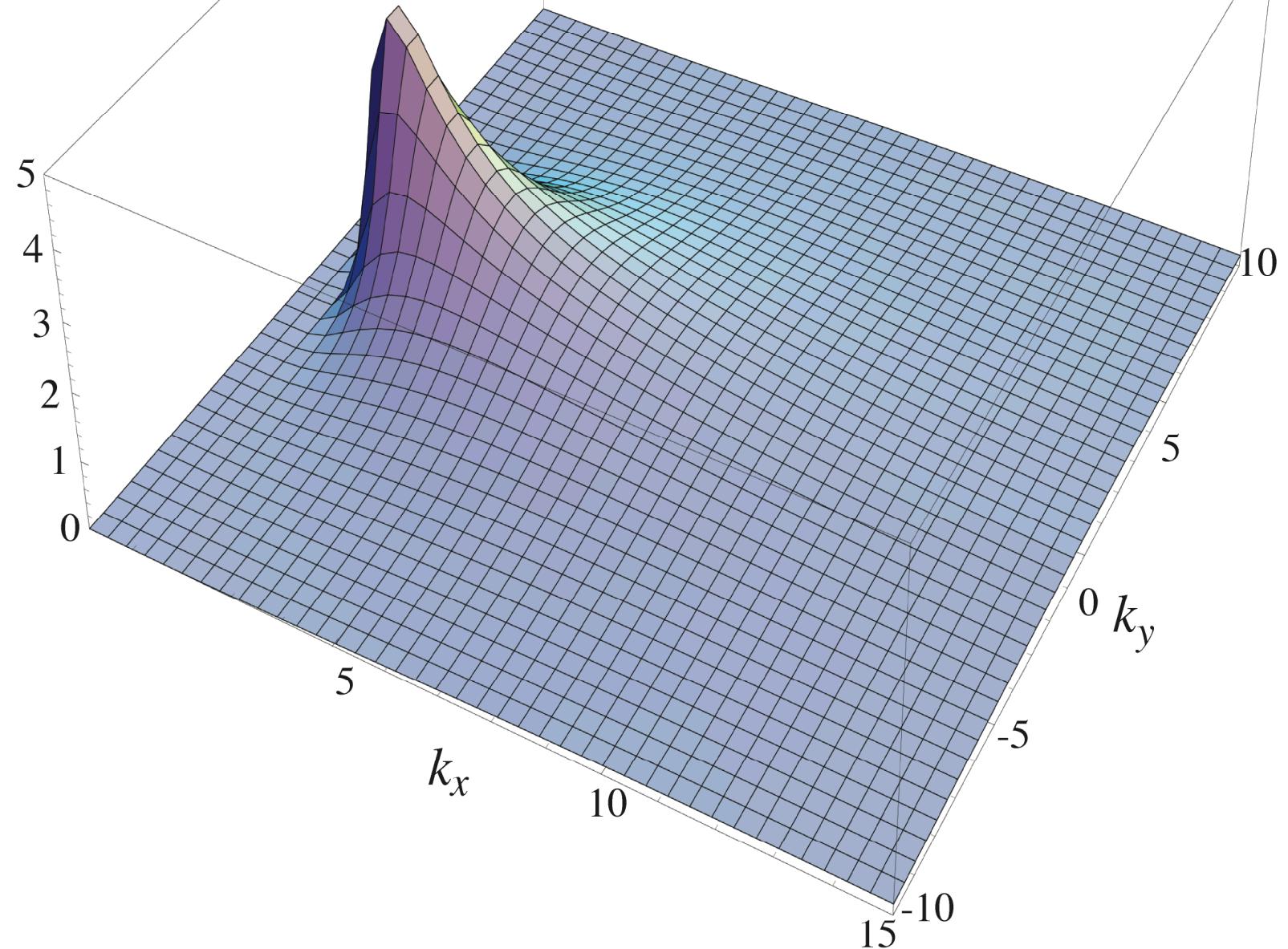
## *Nonlinear Transformation*

$$\eta(x, y, t) = \frac{2}{\lambda} \frac{\partial^2}{\partial x^2} \ln \theta(x, y, t); \quad \lambda = \alpha / 6\beta$$

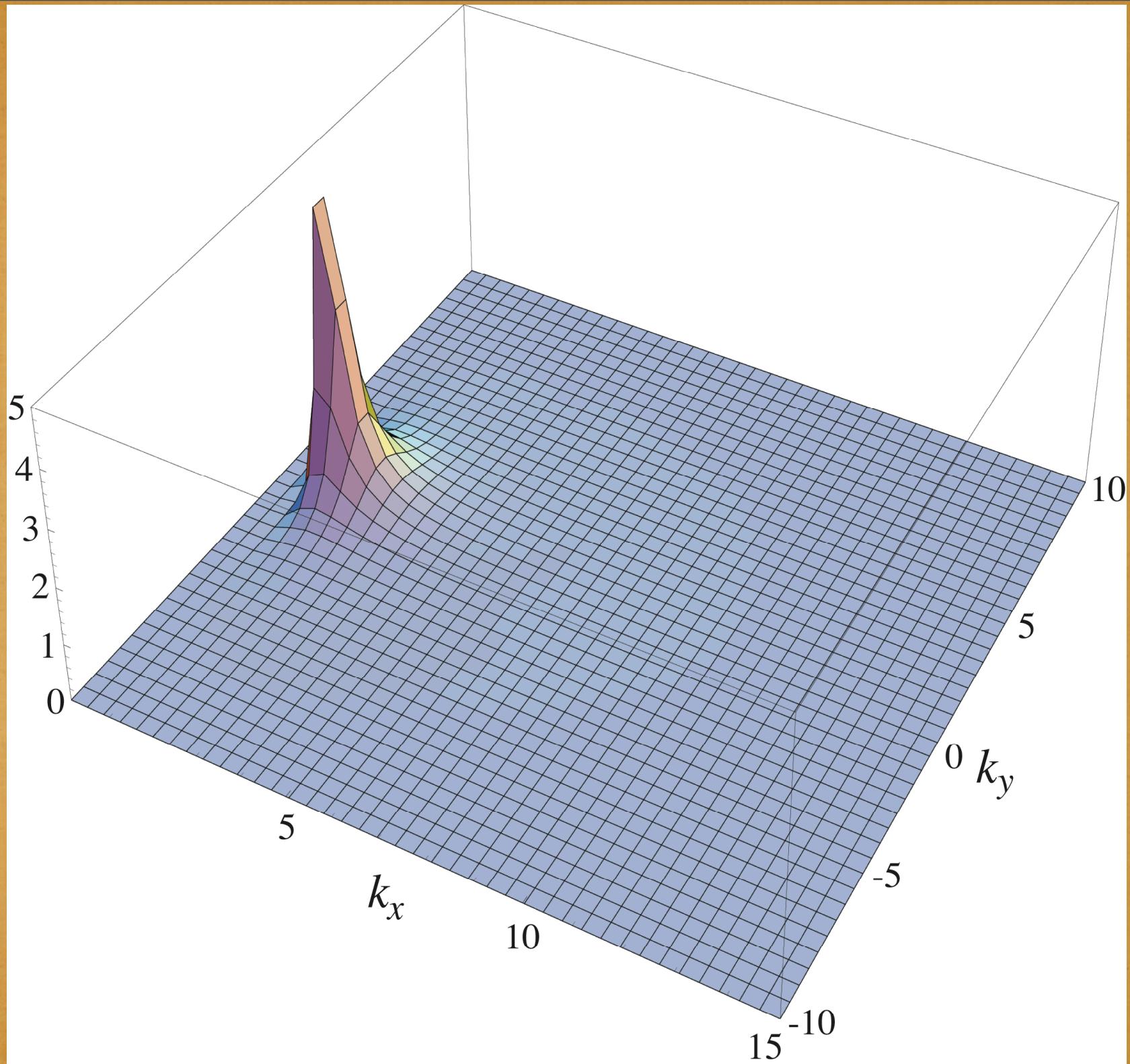
$$\theta(x, y, t) = \exp \left[ \frac{\lambda}{2} \iint_x \eta(x, y, t) dx' dx'' \right]$$

The theta function has a much more narrow spectrum than the surface elevation!

# Surface Elevation spectrum



# Theta Function spectrum



# EULER EQUATIONS

